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Millisecond Pulsars: Detectable Sources of Continuous Gravitational Waves?

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ABSTRACT

Laboratory searches for the detection of gravitational waves have focused on the detection of burst signals emitted during a supernova explosion, but have not resulted in any confirmed detections. An alternative approach has been to search for continuous wave (CW) gravitational radiation from the Crab pulsar. In this paper, we examine the possibility of detecting CW gravitational radiation from pulsars and show that nearby millisecond pulsars are generally much better candidates. We show that the minimum strain $h_c \sim 10^{-26}$ that can be detected by tuning an antenna to the frequency of the millisecond pulsar PSR 1957+20, with presently available detector technology, is orders of magnitude better than what has been accomplished so far by observing the Crab pulsar, and within an order of magnitude of the maximum strain that may be produced by it. In addition, we point out that there is likely to be a population of rapidly rotating neutron stars (not necessarily radio pulsars) in the solar neighborhood whose spindown evolution is driven by gravitational radiation. We argue that the projected sensitivity of modern resonant detectors is sufficient to detect the subset of this population that lies within 0.1 kpc of the sun.

Subject headings: gravitation – radiation mechanisms: nonthermal – stars: neutron – pulsars: general

1. INTRODUCTION

There are two types of signals of gravitational radiation that are expected to be most readily detectable from astrophysical sources: burst signals of short duration, from sources such as a stellar merger or a nonspherical core collapse associated with a supernova event; and continuous wave (CW) signals from sources such as short period, compact binary star systems or rapidly rotating, nonaxisymmetric (or precessing) compact stars. Most experimental searches for gravitational radiation have focused on the detection of burst signals. To date, there have been no confirmed detections of this type of gravitational radiation, although the terminal phase of the

coalescence of neutron-star binaries appears to be a promising source for future ground-based laser interferometers such as LIGO (Abramovici et al. 1992; Cutler et al. 1992; Finn & Chernoff 1993). The detection of CW radiation from binary star systems by a space-based interferometer has been proposed (Faller & Bender 1984; Evans, Iben, & Smarr 1987; Faller et al. 1989). Currently, though, the only experimental effort to detect CW gravitational radiation that is underway has been pioneered by the Tokyo group which has searched for CW emission from the Crab Pulsar (Tsubono 1991). We argue here that nearby millisecond pulsars are likely to be stronger sources of CW radiation than the Crab pulsar and therefore warrant the attention of experimental searches for gravitational radiation.

An axisymmetric object which is rotating about its minor axis will not emit gravitational radiation because it has no time varying quadrupole moment. Therefore a pulsar (or any rotating neutron star, for that matter) must be nonaxisymmetric and/or precessing in order for it to radiate (Ferrari & Ruffini 1969; Zimmerman 1978, 1980; Shapiro & Teukolsky 1983; Barone et al. 1988). Several mechanisms for the production of nonaxisymmetric deformations in pulsars have been suggested (Ipser 1971), including asymmetric crystallization of the crusts (Ruderman 1969; Ferrari & Ruffini 1969), pressure and magnetic stress anisotropies (Ostriker & Gunn 1969, Ruderman 1970), and rotationally induced instabilities (Imamura, Friedman, & Durisen 1985). Misalignments between the symmetry and spin axes of pulsars might occur as the result of electromagnetic torques due to magnetic dipole radiation, corequakes (Pines and Shaham 1974), or encounters with neighboring stars.

Precessing, nearby millisecond pulsars have recently been put forth by de Araújo et al. (1994) as good candidates for the detection of gravitational radiation by upcoming interferometric detectors. They suggest, however, that the signals seen by these detectors due to radiation from wobbling pulsars may be burst signals, not CW signals. This is because the damping time-scale for the wobble angle due to the emission of gravitational radiation is expected to be on the order of seconds whereas the observation time needed to observe CW sources is likely to be $\sim 10^7$ s. This situation does seem likely to present itself with the first generation of detectors because the source signals are expected to be near the limit of detectability. Thus, at least for some time, nonaxisymmetric deformations may be the only channel through which pulsars produce detectable CW radiation.

The rate at which gravitational energy is radiated from a nonaxisymmetric object that is rotating about its minor axis with angular velocity ω is (Ferrari & Ruffini 1969; Shapiro & Teukolsky 1983),

$$\dot{E}_{GR} = -\frac{32G}{5c^5} I_3^2 \epsilon^2 \omega^6. \quad (1)$$

This expression has been derived in the quadrupole approximation for nearly-Newtonian sources assuming that the object has principal moments of inertia I_1 , I_2 and I_3 , respectively, about its three principal axes $a \gtrsim b > c$ fixed in the body frame (Landau & Lifshitz 1962; Misner, Thorne, & Wheeler 1973); and $\epsilon \equiv (a - b)/(ab)^{1/2}$ is the ellipticity in the equatorial plane. Here, and

throughout this paper, dots denote differentiation with respect to time. Several authors have proposed that the Crab pulsar is the best candidate for detecting CW radiation (e.g., Zimmerman & Szedenits 1978; Tsubuno 1991) because it has the largest spin down energy flux density $\dot{E}_{rot}/(4\pi r^2)$ of all pulsars; here r is the distance to the pulsar and $\dot{E}_{rot} = I_3\omega\dot{\omega}$ ($\sim 10^{38}$ erg s $^{-1}$ for the Crab pulsar) is the rotational energy loss rate. We note, however, that for a neutron star of given ellipticity ϵ , the loss rate due to gravitational radiation strongly favors stars that are spinning more rapidly (equation [1]). Also, if the ellipticity is caused by the rapid rotation of the star, one might expect the value of ϵ to be higher in stars that are rotating more rapidly. With this in mind, we have suggested (Barker et al. 1994; see also Schutz 1995) that nearby millisecond pulsars (Backer et al. 1982; Fruchter, Stinebring, & Taylor 1988; Johnston et al. 1993) may be stronger sources of CW radiation than the Crab pulsar (see §2).

One favorable aspect of CW radiation is that a resonant detector can be tuned to the frequency of the emission. In order to tune a cylindrical bar detector so that its primary quadrupole mode of oscillation resonates with the frequency of radiation from the pulsar $\omega_0 = 2\omega$ (see §2), the bar must have a length $L \approx \pi c_s/\omega_0$, where $c_s \approx 5.2 \times 10^5$ cm s $^{-1}$ is the speed of sound in a prototypical aluminum alloy bar (see §3). Thus the length of the bar must be $L \sim 1.3P_{ms}$ m, where $P_{ms} = 10^3(2\pi/\omega)$ is the period of the pulsar in milliseconds. This is prohibitively long ($L \sim 43$ m) for the Crab pulsar which has a period $P_{ms} = 33$. Consequently, the Tokyo group (Tsubuno 1991) has used two short crossed bars instead of a single bar. Because better sensitivities are achievable using a single bar, millisecond pulsars appear to be better candidates from a detector design standpoint as well (see §3).

2. THE CASE FOR MILLISECOND PULSARS

In the quadrupole approximation, the two polarizations of the gravitational strain amplitude h that will be received by an observer located at a distance r along the rotation axis of a source are (Misner, Thorne, & Wheeler 1973; Kochanek et al. 1990; Rasio & Shapiro 1992)

$$h_+ = \frac{G}{c^4} \frac{1}{r} (\ddot{\mathbf{I}}_{xx} - \ddot{\mathbf{I}}_{yy})$$

$$h_\times = \frac{G}{c^4} \frac{2}{r} \ddot{\mathbf{I}}_{xy}. \quad (2)$$

Here \mathbf{I}_{xx} , \mathbf{I}_{yy} , and \mathbf{I}_{xy} are cartesian components of the reduced quadrupole moment,

$$\mathbf{I}_{ij} = I_{ij} - \frac{1}{3} \delta_{ij} \text{Tr} I, \quad (3)$$

where I is the inertia tensor and $\text{Tr} I = (I_1 + I_2 + I_3)$. If the source is a rotating, slightly nonaxisymmetric ($a \sim b, I_1 \neq I_2$) star,

$$\ddot{\mathbf{I}}_{xx} - \ddot{\mathbf{I}}_{yy} = 4\omega^2 I_3 \epsilon \cos(2\omega t)$$

$$\ddot{\mathbf{I}}_{xy} = 2\omega^2 I_3 \epsilon \sin(2\omega t) \quad (4)$$

(Shapiro & Teukolsky 1983) and

$$\begin{aligned} h_+ &= 4 \frac{G \omega^2}{c^4 r} I_3 \epsilon \cos(2\omega t) \\ h_\times &= 4 \frac{G \omega^2}{c^4 r} I_3 \epsilon \sin(2\omega t). \end{aligned} \quad (5)$$

Therefore, the “characteristic” gravitational strain amplitude h_c produced by a star that is spinning down due to the emission of gravitational radiation is

$$h_c \equiv 4 \frac{G \omega^2}{c^4 r} I_3 \epsilon \quad (6a)$$

$$\sim \left(\frac{\dot{E}_{GR}}{1.6 \times 10^{38} \text{ergs}^{-1}} \right)^{1/2} \frac{1}{\omega r}, \quad (6b)$$

where, in the second expression, we have inserted the value of $I_3 \epsilon$ prescribed by equation (1) and the frequency at which the radiation will be detected is $\omega_o = 2\omega$.

2.1. Expected Strains from Pulsars

When a value of $I_3 = 10^{45}$ g cm² is adopted (a reasonable estimate for a $1.4M_\odot$ neutron star of radius 10 km), equation (6a) can be written in the form

$$\frac{h_c}{\epsilon} = 4.2 \times 10^{-18} (P_{ms}^2 r_{kpc})^{-1}, \quad (7)$$

where r_{kpc} is the distance to the star in kiloparsecs. The right-hand side of equation (7) is given entirely in terms of observable pulsar parameters and assumes its largest value ($h_c/\epsilon \sim 10^{-18}$) for nearby millisecond pulsars like PSR 1957+20 and PSR 0437-47 (see Figure 1). Unfortunately, one cannot use this expression to predict a particular strain from a given pulsar without making some assumption about the degree to which the pulsar possesses a nonaxisymmetric structure.

Neutron star ellipticities could conceivably be as large as the breaking strain of their crust, which is thought to be $\lesssim 5 \times 10^{-4}$ (Shapiro & Teukolsky 1983). For any specific pulsar, however, an upper limit to ϵ can be directly inferred from a measurement of \dot{P} if one assumes that the spindown of the pulsar is attributed solely to the emission of gravitational radiation (Press & Thorne 1972). In this situation, $\dot{E}_{rot} = I_3 \omega \dot{\omega}$ can be equated to \dot{E}_{GR} as given in equation (1) and the inferred upper limit to the ellipticity is

$$\epsilon_{GR} = 6.0 (P_{ms}^3 \dot{P})^{1/2}, \quad (8)$$

where, again, we have set $I_3 = 10^{45}$ g cm 2 . Note that, as in equation (7), the right-hand side of this expression is given entirely in terms of observable pulsar parameters. Equation (8) serves only as an upper bound to the ellipticity because other mechanisms are likely to be responsible for some portion of the observed \dot{P} in pulsars. Indeed, in most normal pulsars \dot{P} is thought to be determined primarily by energy losses due to magnetic dipole radiation (Ostriker & Gunn 1969; Pacini 1967) rather than by losses due to gravitational radiation, so the true ellipticities of most pulsars will be less (perhaps orders of magnitude less) than the value ϵ_{GR} given by expression (8). Note that the limit on ϵ set by this expression for millisecond pulsars is very small; for example, for PSR 0437-47 and PSR 1957+20, $\epsilon_{GR} = 3 \times 10^{-8}$ and 1.6×10^{-9} , respectively (see Figure 1 and Table 1).

It proves instructive to construct a plot of the two “observables” h_c/ϵ vs. ϵ_{GR} for a large number of pulsars. Figure 1 includes (as of January, 1995) all of the (~ 470) pulsars in the Caltech Pulsar Database for which P_{ms} , \dot{P} , and r_{kpc} (the distance obtained from the dispersion measurement) were available. It should be noted, first, that because they tend to have very small values of \dot{P} as well as small values of P_{ms} , millisecond pulsars tend to cluster in the upper left-hand corner of this figure. By contrast, normal pulsars tend to cluster in the lower right-hand region of the diagram. The lines drawn in Figure 1, with slopes of -1 , identify the locus of points for sources with identical values of h_{crkpc} as measured from Earth under the (unlikely) assumption that the spindown of all pulsars is attributed solely to the emission of gravitational radiation, i.e., under the assumption $\epsilon = \epsilon_{GR}$. For example, by this assumption, pulsars above and to the right of the solid line would have $h_c > 10^{-26} r_{kpc}^{-1}$ and the Crab and Vela pulsars, in particular, would exhibit the largest values of h_c ($\sim 10^{-24}$).

When confronted with a diagram like Figure 1, one is tempted to assume that the Crab and Vela pulsars will be the brightest CW sources in the sky. However, all that one can safely conclude is that, because these objects exhibit large spindown rates (which are understood to be due to magnetic dipole radiation, not gravitational radiation) the “observable” ϵ_{GR} places only a very loose constraint on the underlying structural ellipticity of these two neutron stars. One is *permitted* by ϵ_{GR} to adopt ellipticity values as large as few $\times 10^{-4}$ — and, hence, strains approaching 10^{-24} — without conflicting with measured values of \dot{P} . However, in light of what we know about millisecond pulsars, it seems unreasonable to assume that these neutron stars have structural ellipticities of this magnitude. In particular, under the premise that nonaxisymmetric structure in neutron stars is likely to be rotationally enhanced (see Imamura, Friedman, & Durisen 1985), one would not expect ϵ to be larger in (the relatively slowly rotating) normal pulsars than it is in millisecond pulsars. Correspondingly, it would be difficult to imagine that the Crab and Vela pulsars have ellipticities larger than $\epsilon \sim 10^{-8} - 10^{-9}$ (the limit set by ϵ_{GR} in millisecond pulsars) and, hence, that they produce strains larger than $h_c \sim 10^{-29} - 10^{-30}$.

In millisecond pulsars, it is conceivable that nonaxisymmetric distortions are induced by the observed, relatively large rotational energies of these objects and that measurable levels of CW radiation result. At first glance, PSR 0437-47 would appear to be the best candidate for

the detection of CW radiation because it is the millisecond pulsar with the largest value of the product $(h_c/\epsilon)\epsilon_{GR} = 2.6 \times 10^{-26}$. However, taking the conservative view that ϵ is nowhere larger than the value of ϵ_{GR} set by PSR 1957+20 (i.e., $\epsilon \leq 1.6 \times 10^{-9}$), the pulsar exhibiting the largest h_c is PSR 1957+20 itself ($h_c = 1.7 \times 10^{-27}$). We consider this to be a much more reasonable estimate of likely strains coming from the best candidates for the detection of CW radiation than the estimates set earlier based on the observed properties of the Crab and Vela pulsars.

2.2. Variation of Characteristic Strain with Age

In models that assume that the spindown of pulsars is due entirely to magnetic dipole radiation, it has been appreciated for quite some time that pulsar evolutionary paths can be drawn in an “observables” diagram similar to our Figure 1 if one adopts a particular function prescribing the strength of pulsar magnetic fields as a function of time (cf., Chanmugam 1992). Similarly, an evolutionary path in the h_c/ϵ vs. ϵ_{GR} diagram can be plotted for any individual pulsar if one assumes that the pulsar’s measured \dot{P} is entirely due to gravitational radiation and one adopts a particular function prescribing the pulsar’s ellipticity as a function of time. To make such an evolutionary discussion relevant to the broad class of pulsars, rather than to one pulsar at a time, we should first remove any distance dependence from the diagram. By analogy with traditional discussions of stellar evolution, we define an “absolute” characteristic strain H_c for any pulsar to be the characteristic strain it would exhibit if it were located a distance $r_{kpc} = 1$ from the sun. Figure 2 re-displays all the pulsar data from Figure 1 on this “absolute” characteristic strain scale.

First, note that if $P_{ms} = 0.5$ is associated with the maximum rotation rate (breakup velocity) of any neutron star, then the dotted horizontal line drawn at $\log_{10}(H_c/\epsilon) = -16.78$ in Figure 2 represents the largest achievable absolute characteristic strain of any neutron star, even at birth. (The true demarcation line may, in fact, be somewhat higher or somewhat lower than this; its correct location cannot be established until our understanding of the equation of state of neutron star matter improves.) The simplest function to choose for $\epsilon_{GR}(t)$ is ϵ_{GR} equals a constant. Then all evolutionary trajectories are vertical in Figure 2 and, because \dot{P} is positive, all trajectories are directed downward. The rate at which a neutron star evolves along its vertical path in Figure 2 is prescribed by equation (8). Specifically, we can write

$$P^3 dP = \left(\frac{\epsilon_{GR}}{A} \right)^2 dt, \quad (9)$$

where $A = 1.9 \times 10^5 \text{ s}^{-3/2}$. Then, assuming for simplicity that the pulsar period initially is much less than its period at time t , the function $P(t)$ becomes

$$P^2 = \left(\frac{2\epsilon_{GR}}{A} \right) t^{1/2}. \quad (10)$$

Combining this with equation (7), the rate of evolution along a (vertical) trajectory in Figure 2 is prescribed by the following expression:

$$\frac{H_c}{\epsilon} = (4.0 \times 10^{-19} s^{1/2}) \epsilon_{GR}^{-1} t^{-1/2}. \quad (11)$$

According to equation (11), a population of neutron stars that have the same age but varying ellipticity will trace out a straight line that has a slope of -1 in Figure 2. More specifically, the solid and dashed line in Figure 2 locate isochrones of $t = 10^8$ and 10^{10} years, respectively. It is interesting to note that none of the individual pulsars plotted in Figure 2 lie below the 10^{10} year, or “Hubble line,” isochrone. Somewhat surprisingly, isochrones in Figure 2 exhibit the same slope (-1) as lines of constant absolute characteristic strain (see Figure 1). Hence, a population of neutron stars of the same age will exhibit the same absolute characteristic strain across the entire population. For reference, the solid and dashed lines in Figure 2 correspond to strains of $H_c = 7.1 \times 10^{-27}$ and 7.1×10^{-28} , respectively.

2.3. A Population of Loud CW Sources

According to models that invoke magnetic dipole radiation to explain the spindown of pulsars, a pulsar’s magnetic field strength is related to P and \dot{P} by the expression $B = 3.2 \times 10^{19} \text{ G s}^{-1/2} (P\dot{P})^{1/2}$ (Chamugam 1992). By this relation, the extremely small value of \dot{P} that has been measured for the millisecond pulsar PSR 1957 + 20 can be understood only if this pulsar possesses a magnetic field that is no stronger than $1.7 \times 10^8 \text{ G}$. It seems unlikely that this field strength is unique among the population of rapidly rotating neutron stars (not all of which are radio pulsars) that resides in the solar neighborhood. Among this sub-population, gravitational radiation will dominate over magnetic dipole radiation as the principal energy loss mechanism throughout each star’s life ($\lesssim 10^{10}$ years) as long as the stars have equatorial ellipticities $\epsilon_{GR} \gtrsim 6 \times 10^{-9}$. (More generally, the criterion works out to be $\epsilon_{GR} \gtrsim 0.20 B_{12}^2$, where B_{12} is the field strength in units of 10^{12} Gauss.)

Given that the tightest constraint that currently has been placed on the ellipticity of neutron stars is $\epsilon \lesssim 5 \times 10^{-4}$, set by the breaking strain of the crust, it does not seem unreasonable to suggest that most neutron stars with $B \lesssim 10^8$ Gauss also have structural ellipticities $\epsilon_{GR} \gtrsim 10^{-8}$. (As a point of reference, according to Cook (1973) the gravitational field of the Earth exhibits an effective ellipticity $\sim 10^{-6}$.) We submit, therefore, that in the solar neighborhood there exists a family of rapidly rotating neutron stars whose spindown is driven by gravitational radiation. According to the preceding discussion, even after 10^{10} years of evolution we are guaranteed that this family of objects will exhibit an absolute strain $\gtrsim 7 \times 10^{-28}$. It would seem, then, that CW detectors designed to reach this strain level are virtually certain to detect this population of stars.

3. EXPERIMENTAL FEASIBILITY

Experiments having the sensitivity necessary to directly detect gravitational radiation are difficult to design and conduct. We consider here a necessary condition for the feasibility of any new experiments designed to improve significantly the sensitivity for detection of CW sources. The specification of *sufficient* conditions is a much broader issue requiring discussion beyond the scope of this paper.

It is well-understood that there are two fundamental, and predictable, noise sources that will necessarily limit the sensitivity of any gravitational wave detector: thermo-mechanical force noise and sensor noise. The first of these arises from random forces applied to the antenna due to its coupling to a finite-temperature heat bath; the second results from random fluctuations at the output of the electro-mechanical devices that are used to sense any gravitationally induced, time-varying distortions of the antenna. For detection of burst sources, both types of noise are important (Solomonson et al. 1993).

The situation is different if one conducts a targeted search for CW emission from a known pulsar. Then information from radio observations can be used to great advantage. For example, from the measured pulsar rotation frequency ω , one knows to what frequency $\omega_o = 2\omega$ the antenna should be tuned; and because \dot{P} is extremely small, one expects to receive a steady flux of gravitational radiation at the prescribed frequency for a very long time ($P/\dot{P} \sim 10^9 - 10^{10}$ yr for millisecond pulsars). The continuous interaction of a resonantly tuned antenna with a steady, periodic gravitational force allows the mechanical amplitude of the antenna to steadily grow with time. Coherent averaging of the motion, accounting for the observer's and source's doppler shifts, can be used to narrow the bandwidth of the measurement to such an extent that the sensor noise, under reasonable conditions, can become completely negligible. Thus it should be possible to avoid sensor noise altogether within a narrow range of frequencies and thereby greatly reduce the minimum detectable strain amplitude h_c compared to the amplitude that is detectable from burst sources. Such techniques have been employed in the search for CW radiation from the Crab pulsar (Tsubono 1991).

In practice, then, the mechanical motion of the antenna will have two parts: the motion induced by the gravitational wave signal and the motion induced by the thermo-mechanical force noise. This remaining noise source is simply the random Langevin force F , or the mechanical equivalent of Johnson-Nyquist noise in electrical circuits. The relative amplitude of these two contributions to the signal can be calculated by starting from the equations of ordinary elastic theory and including the CW gravitational and Langevin forces as driving terms (Merkowitz & Johnson 1994). The elastic equation can be solved via an eigenfunction expansion which produces one harmonic oscillator equation for each eigenmode of the system. In this harmonic oscillator formulation, the CW gravitational force takes the form

$$f_{GW} = \frac{1}{2}\mu L_e \ddot{h}_+(t) = -\frac{1}{2}\mu L_e h_c \omega_o^2 \cos \omega_o t, \quad (12)$$

where μ is the effective mass of the detector and L_e is the effective length of the mode for a particular component of the gravitational strain tensor. Representing the effects of F by a random gravitational field \ddot{h}_R , the Langevin force $f_F(t)$ takes the form

$$f_F(t) = \frac{1}{2}\mu L_e \ddot{h}_R(t). \quad (13)$$

Therefore the total force acting on the antenna can be written as:

$$f_{tot} = \frac{1}{2}\mu L_e(\ddot{h}_R(t) - h_c \omega_o^2 \cos \omega_o t). \quad (14)$$

The spectral density S_R of any random force f_R is given by

$$S_R = \langle \tilde{f}_R^* \tilde{f}_R \rangle, \quad (15)$$

where \tilde{f}_R is the Fourier transform of f_R and an asterisk denotes the complex conjugate. Setting $f_R = f_F$, and realizing that $\ddot{h}_R \cong -\omega_o^2 \tilde{h}_R$, we deduce

$$S_F = \left(\frac{1}{2}\mu L_e \omega_o^2\right)^2 S_h, \quad (16)$$

where

$$S_h \equiv \langle \tilde{h}_R^* \tilde{h}_R \rangle. \quad (17)$$

But the single-sided spectral density of the Langevin force near a resonant mode at frequency ω_o is fixed by simple thermodynamics and the fluctuation-dissipation theorem to be

$$S_F = 4kT \left(\frac{\mu\omega_o}{Q}\right), \quad (18)$$

where k is the Boltzmann constant, T is the physical temperature of the heat bath, and Q is called the mode quality factor (Q^{-1} is the dissipation coefficient). Hence, combining equations (16) and (18) we derive

$$S_h = \frac{16kT}{L_e^2 \mu \omega_o^3 Q}. \quad (19)$$

This expression gives the spectral density of the noise produced in the antenna by the Langevin force, normalized to the amplitude of the force on the antenna due to the gravitational wave itself.

The square root of S_h is referred to as the “strain noise” of the antenna and can be used to define the background noise level below which a signal of astrophysical origin will remain undetectable. More specifically, if one integrates the signal over an observing time τ and demands a 4σ confidence level for detection, the smallest reliably detectable h_c will be

$$h_c \cong 4\sqrt{\frac{S_h}{\tau}}. \quad (20)$$

When developing a detector for CW gravitational radiation, expression (19) is important because it identifies design parameters that contribute most significantly to the amplitude of the

random noise near the resonant mode of the antenna. It is important to realize, in addition, that the factors of length, mass, and frequency are not independent. They are coupled by the geometry and composition of the antenna. For example, for the primary quadrupole mode of a “long” right circular cylinder, or bar, with length L and diameter $d \ll L$, the mode parameters become

$$\mu = \frac{1}{8} \rho \pi d^2 L, \quad (21a)$$

$$L_e \approx \frac{4}{\pi^2} L \approx \frac{4c_s}{\pi\omega_o}, \quad (21b)$$

where ρ is the mass density of the antenna material and c_s is the speed of sound in the material. Substituting these expressions into equation (19), the expression for strain noise becomes

$$\sqrt{S_h} = \left(\frac{8kT}{d^2 \rho c_s^3 Q} \right)^{1/2} = 1.2 \times 10^{-23} \left(\frac{1m}{d} \right) \left(\frac{T}{0.05K} \right)^{1/2} \left(\frac{10^8}{Q} \right)^{1/2} \text{Hz}^{-1/2}, \quad (22)$$

where we have used the density ($\rho \approx 2.7 \text{ g cm}^{-3}$) and sound velocity of aluminum alloy ($c_s \approx 5.2 \times 10^5 \text{ cm s}^{-1}$), which we regard as the most likely material for such an antenna. Note that neither ω_o nor L has a direct effect on the sensitivity of the bar antenna, but the source frequency does affect the length L of the bar ($L = 2.1 \text{ m}$ for PSR 1957+20, for example).

With current technology, the chosen scale parameters are within the realm of feasibility. The specified diameter, for example, is at the limit of current commercial casting capability, and $T \sim 0.05K$ has recently been achieved for the NAUTILUS bar detector under development by the Rome group (Astone 1991). We have assumed that a moderate improvement in the value of Q over presently available bar antennas can be accomplished. Such values of Q have been achieved by annealing small samples of commercially obtained aluminum alloy 5056 (Marsden 1984). The most significant technical difficulty in achieving a strain noise as low as $10^{-23} \text{ Hz}^{-1/2}$, as indicated by our parameterization of expression (22), is likely to be vibration isolation in such an ultra-low-temperature cryostat.

An integration time ~ 4 months seems feasible for tracking a pulsar that is well timed from radio observations, and makes it reasonable to repeat the measurement a number of times for confirmation. Combining expressions (20) and (22), and setting $\tau = 10^7 \text{ s}$, the smallest detectable h_c becomes 1.5×10^{-26} . By comparison, the Tokyo group has obtained an upper limit of $h_c \sim 2 \times 10^{-21}$ for the Crab pulsar (Owa et al. 1988). Current experiments they are conducting are likely to lead to improvements of a factor 100 or so (K. Tsubono, private communication).

A recently proposed, alternative to the cylindrical bar detector is a “spherical” detector having an antenna configured as a truncated icosahedron, that is, an antenna with the same geometric shape as the C_{60} Buckyball molecule (Johnson & Merkowitz 1993; Merkowitz & Johnson 1994). An antenna of this type with a diameter D exhibits 5 resonant quadrupole modes, each having mode parameters of the form:

$$\omega_0 \approx 3.24 \frac{c_s}{D}, \quad (23b)$$

$$\mu = \rho \frac{\pi}{6} D^3, \quad (23a)$$

$$L_e \approx 0.301D. \quad (23c)$$

Again, the size of the antenna is fixed by the sound velocity and the source frequency; for example, $D = 2.2$ m for PSR 1957+20. Substituting these expressions into equation (19) gives the following expression for the strain noise of a spherical antenna:

$$\sqrt{S_h} = \left(\frac{9.92kT}{D^2 \rho c_s^3 Q} \right)^{1/2} = 6.3 \times 10^{-24} \left(\frac{1.6}{P_{ms}} \right) \left(\frac{T}{0.05K} \right)^{1/2} \left(\frac{10^8}{Q} \right)^{1/2} \text{Hz}^{-1/2}. \quad (24)$$

The first equality appears to be little different from the bar case but, actually, because the diameter D for a sphere is larger than the diameter d for a bar, when both have the same resonant frequency the value of $\sqrt{S_h}$ for the sphere is smaller. Additional advantages of a spherical detector are that it can be equally sensitive to a wave from any direction and it is capable of measuring the direction and polarization of the wave (Forward 1971; Wagoner & Paik 1976; Merkowitz & Johnson 1994).

In the second equality of expression (24) we have used the coupling of the antenna's diameter to the frequency to eliminate D in favor of the pulsar period P_{ms} . We have again assumed aluminum alloy as the material. As in the case of the bar detector, a significant technical difficulty will be maintaining high Q when casting or joining aluminum to make a sphere of diameter 2.2 m designed to resonate with twice the rotation frequency of PSR 1957+20. If the integration time is again chosen to be $\tau = 10^7$ s, the minimum detectable strain for a spherical detector tuned to PSR 1957+20 is predicted to be $h_c = 8.0 \times 10^{-27}$.

4. DISCUSSION

Nearby millisecond pulsars are good candidates for the detection of CW gravitational radiation. Because of their close proximity and rapid rotation, they are capable of emitting radiation with larger gravitational strain amplitudes than pulsars that are farther away and/or have longer periods.

The minimum strain $h_c \sim 10^{-26}$ (1.5×10^{-26} for a bar detector, 8.0×10^{-27} for a spherical detector) that can be detected by designing an antenna tuned to the rotation frequency of the millisecond pulsar PSR 1957+20 and employing presently available resonant detector technology is orders of magnitude better than what has been accomplished so far by observing the Crab pulsar, and within an order of magnitude of the maximum strain that can be produced by PSR 1957+20 as a result of rotationally induced nonaxisymmetric deformations. The design and operation of a resonant antenna that is tuned to the rotation frequency of PSR 1957+20 would, at the very least, place physically meaningful constraints on the nonaxisymmetric ellipticity of millisecond pulsars.

We have argued (§2.3) that there almost certainly is a population of rapidly rotating neutron stars (not necessarily radio pulsars) within the solar neighborhood whose spindown evolution is driven by gravitational radiation. Throughout their entire lifetime, these stars will radiate at an “absolute” strain $H_c \gtrsim 7 \times 10^{-28}$. It is significant that the projected sensitivity of modern resonant detectors is sufficient to detect the subset of this population of stars that resides within 0.1 kpc of the sun.

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Table 1: Data for Representative Pulsars^a

Pulsar	P _{ms}	\dot{P}	r _{kpc}	h _c / ϵ	ϵ_{GR}
Crab	33.4	4.21E-13	2.49	1.52E-21	7.55E-4
Vela	89.3	1.25E-13	0.61	8.70E-22	1.80E-3
PSR 0437-47	5.76	1.2 E-19	0.14	9.12E-19	2.89E-8
PSR 1957+20	1.61	1.68E-20	1.53	1.07E-18	1.59E-9

^aFrom the Caltech Pulsar Database

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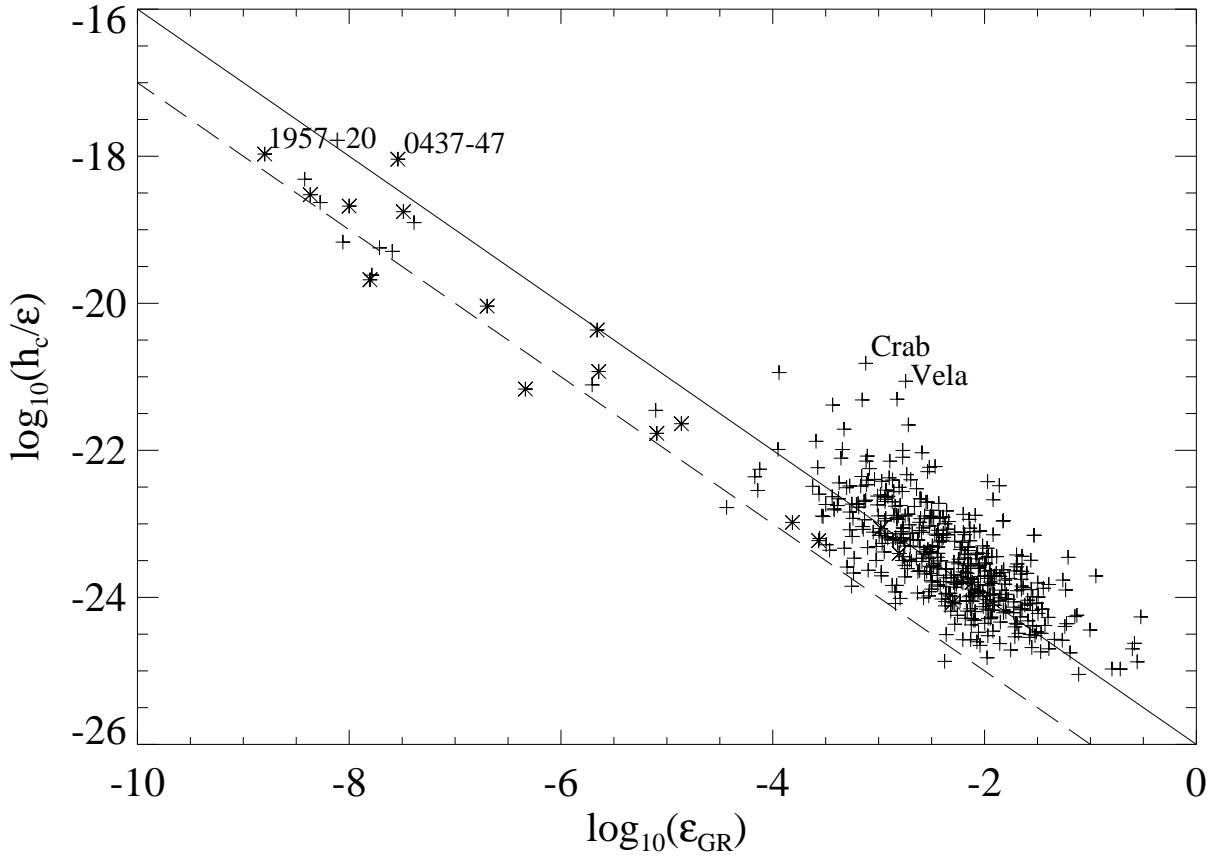


Fig. 1.— The two “observables” $h_c/\epsilon = 4.2 \times 10^{-18} (P_{ms}^2 r_{kpc})^{-1}$ versus $\epsilon_{GR} = 6.0 (P_{ms}^3 \dot{P})^{1/2}$, for ~ 470 pulsars from the Caltech Pulsar Database. Binary pulsars are denoted by asterisks. The solid and dashed lines identify the locus of points for sources with identical values of $h_{cr_{kpc}} = 10^{-26}$ and 10^{-27} , respectively, under the assumption that gravitational radiation is solely responsible for the spindown of the pulsars, and therefore, that $\epsilon = \epsilon_{GR}$.

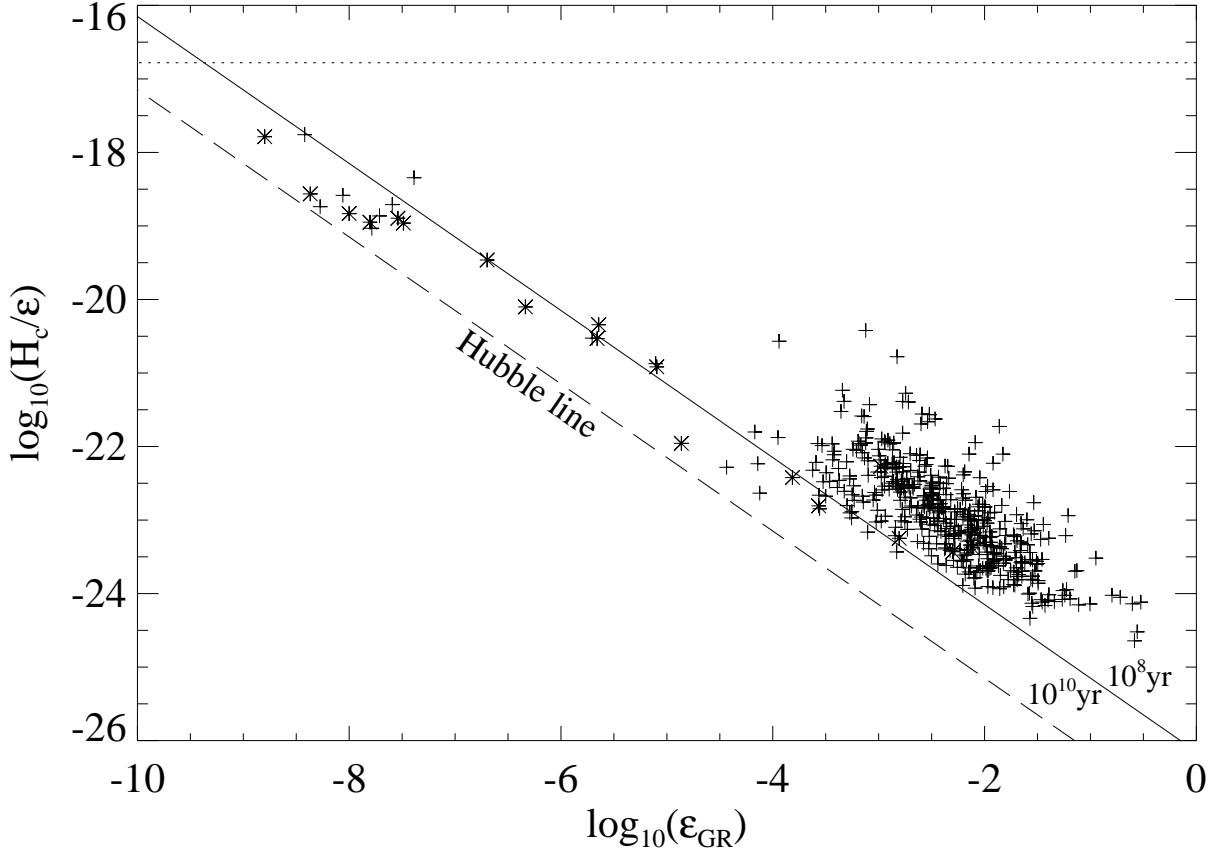


Fig. 2.— The same data shown in Figure 1, except that all pulsars are now assumed to be located at a distance $r_{kpc} = 1$ from the sun. On this scale, the “absolute” characteristic gravitational strain amplitude is $H_c \equiv 4.2 \times 10^{-18} P_{ms}^{-2} \epsilon$. If the breakup velocity limits the period of a neutron star to $P_{ms} = 0.5$, the dotted horizontal line marks the largest possible H_c/ϵ for any neutron star. The solid and dashed lines represent the isochrones of pulsars with ages of 10^8 and 10^{10} (the Hubble line) yrs, respectively, under the assumptions that ϵ_{GR} is a constant function of time and $\epsilon = \epsilon_{GR}$.